Finite Element Methods: Numerical Exercises

Paul J. Atzberger

- 1. Show that each of the elements have the stated regularity as follows:
 - (a) Lagrange triangular element based on \mathcal{P}_k with k+1 distinct nodes along each edge is C^0 .
 - (b) Hermite triangular element based on \mathcal{P}_3 is C^0 .
 - (c) Argyris triangular element based on \mathcal{P}_5 is \mathbb{C}^1 in the normal direction across edges.

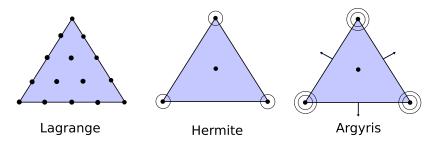


Figure 1: Triangular Elements.

2. There are many ways to develop quadratures for triangulations \mathcal{T} to approximate

$$\int \int_{\mathcal{T}_0} f(\mathbf{x}) d\mathbf{x} \approx \sum_k w_k f(\mathbf{x}_k), \quad \mathbf{x} = (x_1, x_2).$$

(a) Consider Duffy's Transform from a reference triangular element to a quadrilateral element as shown in Figure 2. This is given by

$$\xi = \left(\frac{1+\xi'}{2}\right) \left(\frac{1-\eta'}{2}\right), \quad \eta = \frac{1+\eta'}{2}$$
 $\xi' = \frac{2\xi}{1-\eta} - 1, \quad \eta' = 2\eta - 1,$

where $\eta \in [0,1]$, $\xi \in [0,1-\eta]$, $\xi', \eta' \in [-1,1]$. We can express integration over the triangular element as

$$\int_0^1 \int_0^{1-\eta} f(\xi, \eta) d\xi d\eta = \int_{-1}^1 \int_{-1}^1 f(\xi, \eta) J(\xi', \eta') d\xi' d\eta',$$

where the Jacobian for Duffy's Transform is given by $J(\xi', \eta') = \frac{1}{8}(1 - \eta')$.

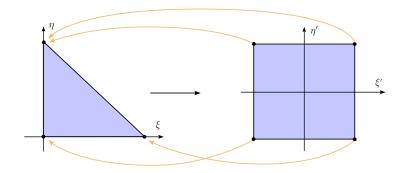


Figure 2: Duffy's Transform.

Use Gaussian quadratures for the cases of 2 and 3 nodes to construct quadratures for the iterated integrals for the quadrilateral. Determine the corresponding nodes and weights for the triangle and construct a quadrature table for the triangular elements for the Gaussian quadrature cases.

(b) Alternatively, we can use for the weights w_k and nodes \mathbf{x}_k from Table 1. For n=4,7, compare this with the Duffy's Transform approach for the test functions (i) $3x^3y^2$, (ii) $\sin(\pi xy/2)$, and (iii) $\exp(-3x^2+3y^2)$. In each case, which yields the more accurate approximation.

d	n	k	\mathbf{x}_k	ω_k	k	\mathbf{x}_k	w_k	k	\mathbf{x}_k	w_k	k	\mathbf{x}_k	w_k
1	1	1	(1/3,1/3)	1/2									
2	3	1	(1/6,1/6)	1/6	2	(2/3,1/6)	1/6	3	(1/6,2/3)	1/6			
3	4	1	(1/3,1/3)	-9/32	2	(3/5,1/5)	25/96	3	(1/5,3/5)	25/96	4	(1/5,1/5)	25/96
4	7	1	(0,0)	1/40	2	(1/2,0)	1/15	3	(1,0)	1/40			
		4	(1/2,1/2)	1/15	5	(0,1)	1/40	6	(0,1/2)	1/15	7	(1/3,1/3)	9/40

Table 1: Quadratures on triangulations for $\int_0^1 \int_0^{1-x_1} f(\mathbf{x}) d\mathbf{x} \approx \sum_k f(\mathbf{x}_k) w_k$, $\mathbf{x} = (x_1, x_2)$. The d is the quadrature order, n number of nodes, \mathbf{x}_k nodes, and ω_k weights. For affine reference element map $\mathbf{x} = \psi(\mathbf{X})$ with $\psi(\mathcal{T}_\ell) = \mathcal{T}_0$ and Jacobian $J(\mathbf{X}) = |\det \partial \psi / \partial \mathbf{X}|$, the quadrature is applied using $\int_{\mathcal{T}_\ell} F(\mathbf{X}) d\mathbf{X} = \int_{\mathcal{T}_0} F(\psi^{-1}(\mathbf{x})) J^{-1} d\mathbf{x}$.

3. Consider the elliptic PDE (Poisson problem) given by

$$\Delta u(\mathbf{x}) = -f(\mathbf{x}), \ \mathbf{x} \in \Omega, \ u(\mathbf{x}) = g(x), \ \mathbf{x} \in \partial \Omega,$$

where $\Omega = [-L, L] \times [-L, L] \subset \mathbb{R}^2$, and g(x) are the boundary values. In the Ritz-Galerkin approximation, we seek a solution $u_h \in \mathcal{V}_h \subset \mathcal{V} = H_0^1(\Omega)$ with

$$a(u_h, w) = \langle f, w \rangle_{L^2}, \ \forall w \in \mathcal{V}_h,$$

where $a(u_h, w) = \int_{\Omega} \nabla_{\mathbf{x}} u_h(\mathbf{x}) \cdot \nabla_{\mathbf{x}} w(\mathbf{x}) d\mathbf{x}$ and $\langle f, w \rangle_{L^2} = \int_{\Omega} f(\mathbf{x}) w(\mathbf{x}) d\mathbf{x}$. Consider a basis of functions $\{\phi_k\}_{k=1}^N$ for \mathcal{V}_h . We can represent any $v \in V_h$ by $v(\mathbf{x}) = \sum_i v_i \phi_i(\mathbf{x})$,

 $u_h(\mathbf{x}) = \sum_i u_i \phi_i(\mathbf{x})$, and approximate f by $f_h(\mathbf{x}) = \sum_i f_i \phi_i(\mathbf{x})$. The FEM approximation u_h can be expressed as solving the linear system

$$A\mathbf{u} = M\mathbf{f}$$
.

The A is the stiffness matrix given by $A_{ij} = a(\phi_i, \phi_j)$, M is the mass matrix given by $M_{ij} = \langle \phi_i, \phi_j \rangle_{L^2}$, and $[\mathbf{u}]_i = u_i$, $[\mathbf{f}]_i = f_i$.

To handle the Dirichlet boundary conditions we need to use that the boundary values g(x) determine some of the nodal variables. By ordering the nodal indices appropriately, we can split the system into components as $\mathbf{u} = [\mathbf{u}_I, \mathbf{u}_B]$ and $A = [A_I|A_B]$. The \mathbf{u}_I corresponds to the nodal locations interior to the domain Ω and \mathbf{u}_B correspond to the nodal locations on the boundary $\partial\Omega$. Since the values \mathbf{u}_B are known, be sure to move these to the right-hand-side (RHS) of the linear system when solving. By restricting to the rows of the system for the indices of \mathbf{u}_I , we obtain the linear system $A_I\mathbf{u}_I = M\mathbf{f} - A_B\mathbf{u}_B$.

(a) (Meshing) Discretize the domain Ω into elements $\mathcal{T} = \{\mathcal{T}_{\ell}\}_{\ell=1}^{m}$, where \mathcal{T}_{ℓ} are triangular elements. For the square domain $\Omega = [-L, L] \times [-L, L] \subset \mathbb{R}^{2}$, one way to discretize is to define a coarse mesh. A basic algorithm to obtain a more refined discretization is to loop over each triangle and bisect the edges to obtain four smaller triangles, see Figure 3. Data structures for this are a list of vertices $\mathbf{v}_{i} \in \mathbb{R}^{2}$ and tuples (i_{1}, i_{2}, i_{3}) which give the indices of the vertices of each triangle.

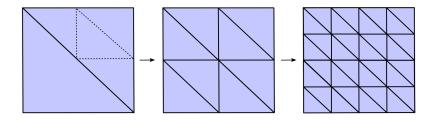


Figure 3: Mesh triangulation and refinement by triangle bisection.

Implement this meshing algorithm for the triangulation in Figure 3. Plot the triangulations when this refinement procedure is done up to n = 5 times.

(b) (Assembly and Quadratures) For the discretization into triangular elements $\mathcal{T} = \{\mathcal{T}_{\ell}\}_{\ell=1}^m$, take $\{\phi_k\}_{k=1}^N$ to be the nodal basis functions for Lagrange elements with polynomial shape functions of degree d so that $v_h | \mathcal{T}_{\ell} \in \mathcal{P}_d$. The stiffness matrix A is obtained through an assembly procedure where we compute the integral by breaking it into parts summing up the inner-products over each element \mathcal{T}_{ℓ} as $A_{ij} = a(\phi_i, \phi_j) = \sum_{\ell=1}^m \int_{\mathcal{T}_{\ell}} \nabla_{\mathbf{x}} \phi_i(\mathbf{x}) \cdot \nabla_{\mathbf{x}} \phi_j(\mathbf{x}) d\mathbf{x} = \sum_{\ell=1}^m A_{\ell,ij}$, and similarly, $M_{ij} = \langle \phi_i, \phi_j \rangle_{L^2} = \sum_{\ell=1}^m \int_{\mathcal{T}_{\ell}} \phi_i(\mathbf{x}) \phi_j(\mathbf{x}) d\mathbf{x} = \sum_{\ell=1}^m M_{\ell,ij}$. Integrals are approximated by high-precision quadratures

$$\tilde{A}_{\ell,ij} = \sum_{k} \omega_k \nabla_{\mathbf{x}} \phi_i(\mathbf{x}_k) \cdot \nabla_{\mathbf{x}} \phi_j(\mathbf{x}_k), \quad \tilde{M}_{\ell,ij} = \sum_{k} \omega_k \phi_i(\mathbf{x}_k) \phi_j(\mathbf{x}_k).$$

The $\{\omega_k\}$ are the quadrature weights and $\{\mathbf{x}_k\}$ are the quadrature nodes. Note in general the quadrature nodes can differ from the finite element nodes. We use these approximations to obtain

$$\tilde{A}_{11} = \tilde{M}_{f}$$
.

For the case of Lagrange elements using polynomial spaces of degree d, we use quadratures that have order 2d. This allows for computing the integrals up to round-off errors. For quadratures on triangulations, see Figure 4 and Table 1.

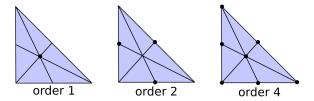


Figure 4: Quadrature Nodes.

Using this assembly + quadrature approach, implement codes to compute for a given triangulation the stiffness and mass matrices when d = 1 and d = 2.

Consider the FEM approximation for the solutions u with $L = \pi$ and (i) $u(x_1, x_2) = \cos(5x_1)\sin(5x_2)$ and (ii) $u(x_1, x_2) = \exp(-\cos(3x_1) + \sin(3x_2))$. Use $f(\mathbf{x}) = -\Delta u$ evaluated at the nodal points to obtain the numerical data for these test problems.

Make a log-log plot of the solution error vs mesh size $h^{-1} = 2^{-n}$ for meshes with refinements n = 1, 2, ..., 5. What is the exhibited order of accuracy of the Lagrange FEMs when d = 1 and d = 2?

(c) (Iterative Methods) To solve approximately

$$A\mathbf{u} = \mathbf{b}$$
, where $\mathbf{b} = M\mathbf{f}$,

iterative methods can be used of the form

$$B\mathbf{u}^{n+1} = C\mathbf{u}^n + \mathbf{b}.$$

For convergence, B-C=A and the spectral radius of $B^{-1}C$ is taken to satisfy $\rho(B^{-1}C)<1$. It is common to decompose the matrix as A=D-L-U, where D is the diagonal entries, -L the lower entries, and -U the upper entries. A few example iterative methods are

- i. Direct Relaxation with B = I and $C = I + \eta A$, with small enough η s.t. $\eta \le 2/\lambda$ or smaller, where λ is the largest eigenvalue of A.
- ii. Jacobi Iteration with B = D and C = L + U.
- iii. Gauss-Seidel Iteration with B = D + L and C = U.

Compare these methods for approximating the solution **u** when $L = \pi$ and (i) $u(x_1, x_2) = \cos(5x_1)\sin(5x_2)$ and

(ii) $u(x_1, x_2) = \exp(-\cos(3x_1) + \sin(3x_2))$. Use $f(\mathbf{x}) = -\Delta u$ evaluated at the nodal points to obtain the numerical data for these test problems.

Make a log-log plot of the number iterations and the error for meshes with n=5 refinements. How many iterations does each method need to converge to 1% accuracy for solving the linear system? We remark that in practice these convergence rates are further enhanced by using preconditioners.